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# **STARK BROADENING CALCULATIONS OF NEUTRAL COPPER SPECTRAL LINES AND TEMPERATURE DEPENDENCE**

**B. Zmerli, N. Ben Nessib, M. S.  
Dimitrijević and S. Sahal Bréchet**

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# **Introduction**

**In plasmas, Stark broadening of spectral lines is important not only for experimental methods, but also for theoretical understanding.**

**The neutral copper (Cu) is often used in electrical industry as an electrode material. Therefore, diagnostic techniques for this element are of particular interest in industrial laboratories.**

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## **Stark broadening calculation:**

1. semi-classical method
2. Impact approximation

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# SC method

S. Sahal-Bréchet, *Astron. Astrophys.* **1**, 91  
(1969)

S. Sahal-Bréchet, *Astron. Astrophys.* **2**, 322  
(1969)

M.S. Dimitrijević, S. Sahal-Bréchet, *Phys. Rev.*  
**A 31**, 316 (1985)

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*In 2008, in order to demonstrate the deviations from the  $T^{-1/2}$  dependence for neutral atom lines, we derived the explicit temperature dependence for two particular cases: the simplified formulae of Freudenstein and Cooper and Dimitrijević and Konjević :*

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## Temperature dependence of atomic spectral line widths in a plasma

B. Zmerli<sup>1</sup>, N. Ben Nessib<sup>1</sup>, and M.S. Dimitrijević<sup>2,\*</sup>

<sup>1</sup> Groupe de Recherche en Physique Atomique et Astrophysique, Institut National des Sciences Appliquées et de Technologie, Centre Urbain Nord, BP 676, 1080 Tunis Cedex, Tunisia

<sup>2</sup> Astronomical Observatory, Volgina 7, 11160 Belgrade, Serbia

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**Abstract.** We investigated here temperature dependence of Stark widths for neutral atom spectral lines in order to find a more precise method for scaling with temperature than sometimes used dependence  $T^{-1/2}$ ,

In this work, we present semiclassical calculations of Stark broadening ( $W$  and  $d$  en Å) for neutral copper spectral lines:

$$w = N \int_0^{\infty} \nu f(\nu) d\nu \left( \sum_{i' \neq i} \sigma_{ii'}(\nu) + \sum_{f' \neq f} \sigma_{ff'}(\nu) + \sigma_{el} \right),$$

$$\sum_{i' \neq i} \sigma_{ii'}(\nu) = \frac{1}{2} \pi R_1^2 + \int_{R_1}^{R_2} 2 \pi \rho d\rho \sum_{i' \neq i} P_{ii'}(\rho, \nu),$$

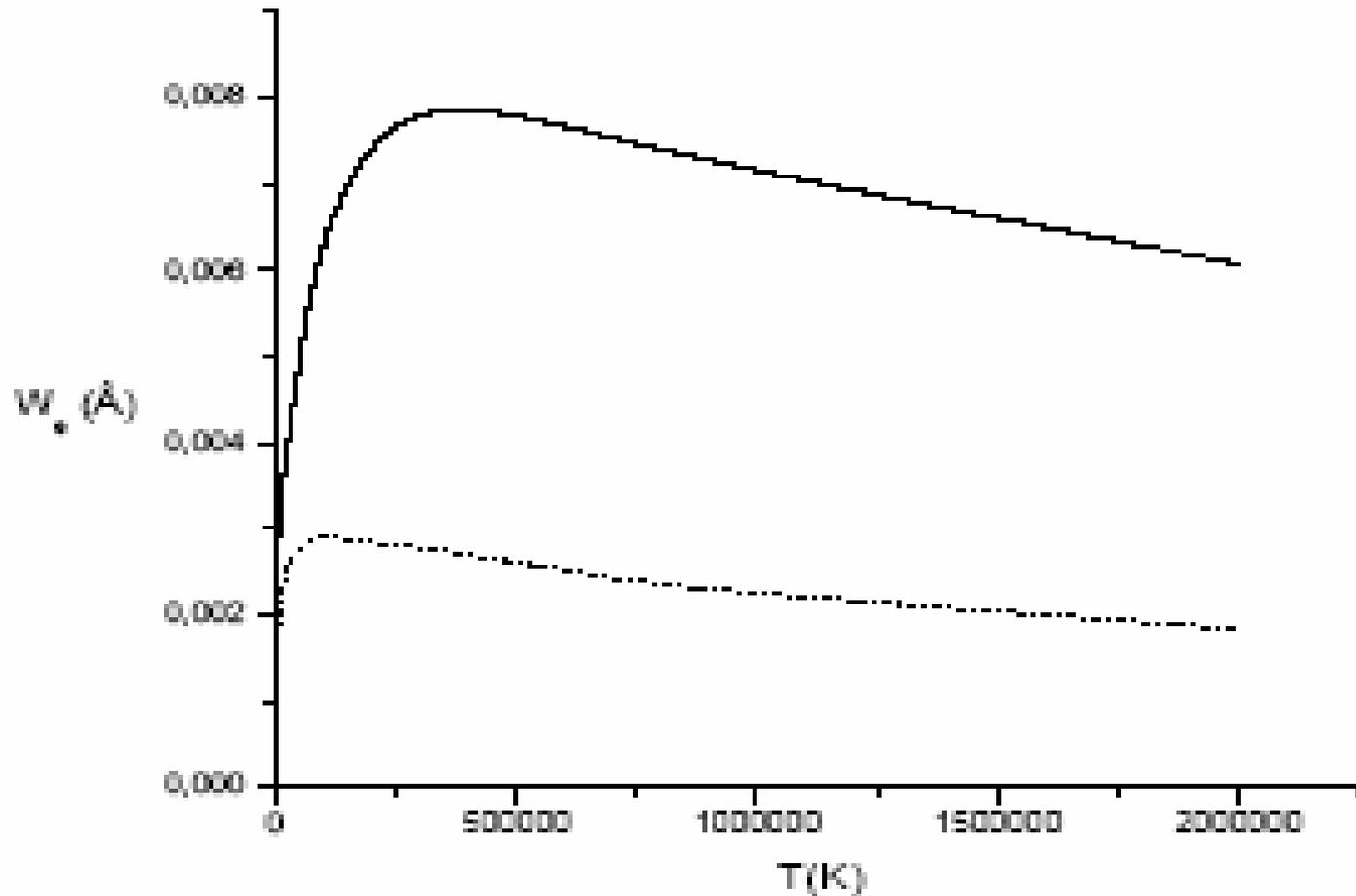
$$\sigma_{el} = 2 \pi R_2^2 + \int_{R_1}^{R_2} 8 \pi \rho d\rho \sin^2 \delta, \delta = (\phi_p^2 + \phi_q^2)^{1/2}.$$

$$d = N \int_0^{\infty} \nu f(\nu) d\nu \int_{R_3}^{R_D} 2 \pi \rho d\rho \sin^2 \phi_p.$$

$$A = \left( \frac{e F_0^2}{\hbar W_e} \left| \alpha_i - \alpha_f \right| \right)^{\frac{3}{4}},$$

$$F_0 = 2 \pi \left( \frac{4}{15} \right)^{\frac{2}{3}} e N^{\frac{2}{3}}, \alpha_i = 4 a_0^3 \sum_{i' \neq i} f_{ii'} \left( \frac{I_H}{\Delta E_{ii'}} \right)^2.$$

In the figure, we show Stark width for the copper resonance spectral line (3247.54 Å): dot line – electron width for 3-level model and solid line - electron width for multi-level model. We calculate the critical temperature  $T_0$  (=361 kK) corresponding to the maximal value of the width. The width increases before  $T_0$  then it decreases.



$$T_0 = \left( \frac{R_{jj'} \Delta E_{jj'}}{3 k_B} \right),$$

$$R_{jj'}^2 = f_{jj'} \frac{\lambda_{jj'} (\text{\AA})}{303.7}$$

$\lambda$ (Å)	Transition	$T_0$ (kK)
5105.54	$4s^2 \ ^2D[5/2] - 4p \ ^2P^o[3/2]$	292
5700.24	$4s^2 \ ^2D[3/2] - 4p \ ^2P^o[3/2]$	370
5782.13	$4s^2 \ ^2D[3/2] - 4p \ ^2P^o[1/2]$	301
3273.96	$4s \ ^2S[1/2] - 4p \ ^2P^o[1/2]$	398
3247.54	$4s \ ^2S[1/2] - 4p \ ^2P^o[3/2]$	361

# DISCUSSION

An improved agreement is found between our results and the theoretical ones.

The ratio  $(W_e/W_1)$  is about 78%,  $(W_e/W_2) = 76\%$ ,  $(W_e/W_4) = 51\%$  and the agreement with Babina results  $W_3$  reaches 95%.

$(d_e/d_1) = 84\%$ ,  $(d_e/d_2) = 85\%$  and  $(d_e/d_3) = 42\%$ .

$(A/A_1) = 105\%$ ,  $(A/A_2) = 111\%$  and  $(A/A_3) = 90\%$ .

For experimental results the agreement is acceptable:

$W/W_K = 60\%$ ,  $W/W_L = 32\%$ ,  $W/W_F = 26\%$ ,  $W/W_S = 33\%$  and  $d/L = 20\%$ .

We have a good agreement with Fleurier result  $d/d_F = 151\%$ .

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# CONCLUSION

This method of Stark width calculation can be widely used in getting more accurate results were extensive set of data are needed.

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**Thank you**

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